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Fluorescence-excitation and Emission Spectroscopy on Single FMO Complexes

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In green-sulfur bacteria sunlight is absorbed by antenna structures termed chlorosomes, and transferred to the RC via the Fenna-Matthews-Olson (FMO) complex. FMO consists of three monomers arranged in C_3 symmetry where each monomer accommodates eight Bacteriochlorophyll *a* (BChl *a*) molecules. It was the first pigment-protein complex for which the structure has been determined with high resolution and since then this complex has been the subject of numerous studies both experimentally and theoretically. Here we report about fluorescence-excitation spectroscopy as well as emission spectroscopy from individual FMO complexes at low temperatures. The individual FMO complexes are subjected to very fast spectral fluctuations smearing out any possible different information from the ensemble data that were recorded under the same experimental conditions. In other words, on the time scales that are experimentally accessible by single-molecule techniques, the FMO complex exhibits ergodic behaviour.

Photosynthesis is one of the most important energy conversion processes on Earth. The generic principle of photosynthesis is that sunlight is absorbed by antenna complexes and then this energy is transferred efficiently to the photochemical reaction centre that acts as a transducer. Typically, both the antennae and the reaction centre (RC) are pigment-protein complexes accommodating a number of pigments that are properly positioned with respect to each other by the interaction with a protein scaffold. For many of these complexes the structural arrangement has been elucidated with high resolution^{1–9}, stimulating enormous interest in the investigation of the structure-function relationships of these systems¹⁰. Next to ultra-fast time-resolved^{11–15} and high-resolution spectroscopy^{16–23} single-molecule techniques in particular have contributed tremendously to our current understanding of the electronic structure of those pigment-protein complexes^{24–30}.

In green-sulfur bacteria sunlight is absorbed by antenna structures termed chlorosomes (a unique case where no proteins are involved) and transferred to the RC via the Fenna-Matthews-Olson (FMO) complex, which is thought to play an important role in unidirectional energy transfer³¹. As a matter of fact the FMO complexes of the bacterial species *Chlorobium limicola* and *Prosthecochloris aestuarii* were the first photosynthetic antenna complexes for which a high-resolution x-ray structure was determined^{32,33}. Meanwhile the structure of the FMO complex from the species *Chlorobaculum (Cba.) tepidum* also became available³⁴, and resembles closely the former one. Accordingly, this FMO complex consists of three identical subunits that are arranged in C_3 symmetry with respect to each other, such that the C_3 axis is perpendicular to the membrane axis, Fig. 1a. Each subunit contains 7 closely interacting bacteriochlorophyll (BChl) *a* molecules and an additional 8th BChl *a* pigment that is located in a region directed towards the chlorosome³⁵, Fig. 1b,c. Numerous studies including linear spectroscopy, linear and circular dichroism, hole burning, transient absorption, and photon echo spectroscopy flanked by theoretical work and computer simulations have been performed on this system (reviewed in ref. 31), making this complex one of the most widely studied pigment-protein complexes so far. Unfortunately, the site energies of the pigments cannot be deduced directly from the ensemble absorption spectrum which features rather broad bands. As a consequence of this, the available information about the excitonic coupling between the BChl *a* pigments is based on theoretical calculations and/or numerical fits of the optical spectra, and the outcome is sensitive to the choice of the input parameters. Although the models discussed differ slightly in their details there is agreement that the electronic coupling between the seven BChl *a* in close proximity leads to the formation of exciton states

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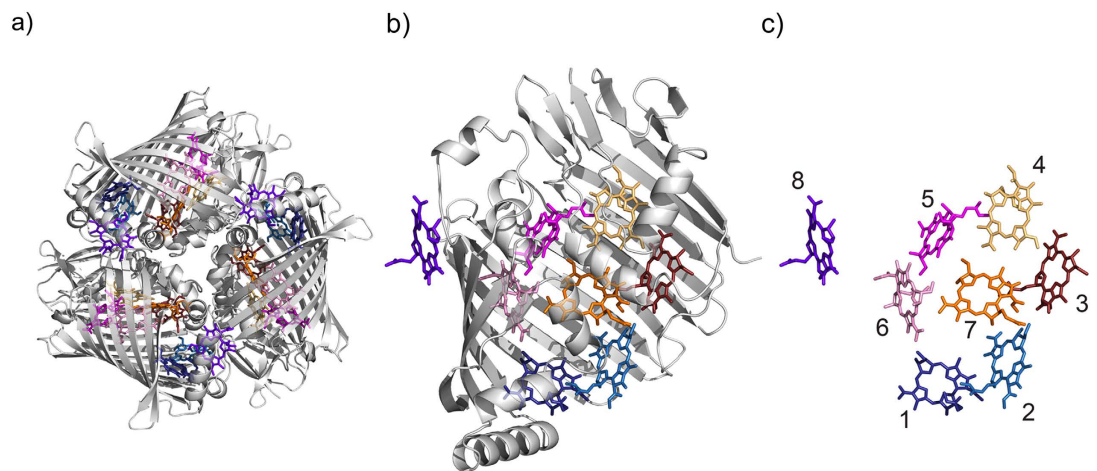


Figure 1. Structure of the FMO complex. (a) Top-view of the trimeric FMO complex from *Chlorobaculum tepidum*. (b) Side view of the BChl *a* arrangement within one monomer unit. For clarity the phytol tails have been omitted. (c) Same as (b) without the protein scaffold. The delocalisation of the excitation over two pigments is indicated by similar colour tones (violet, blue, orange). Structural data taken from www.rcsb.org (protein code 3ENI).

that are delocalised over only a few monomers³⁶. Upon excitation the decay from the highest to the lowest exciton state occurs on a (sub-)picosecond time scale, and the population can be transferred either by a few big steps or by many small steps involving all intermediate exciton states. Finally, the excitation energy leaves the FMO complex en route to the RC via BChl *a* number “3” (“exit pigment”), see Fig. 1b,c.

Recently, by applying 2D spectroscopy, long-lived quantum coherences among the electronic excitations of the pigments in the FMO complex have been observed^{37,38}, and it has been argued that these reflect correlations of the site energy fluctuations. A technique that might shed some light on this debate is single-molecule spectroscopy, because, in general, the broadening of the spectra due to averaging over sample heterogeneities can be avoided. Indeed, there is a demand for experiments on single FMO complexes in the community and challenging experimental schemes have been proposed recently³⁹. Here we present the fluorescence-excitation and emission spectra from individual FMO complexes from the species *Cba. tepidum* and find that these complexes behave ergodically on the experimental accessible timescales.

Results

In order to select a single well-separated FMO for fluorescence-excitation spectroscopy we recorded a wide-field fluorescence image from the sample. An example for such an image is shown in Fig. 2. It features a few diffraction-limited bright spots, each corresponding to an individual FMO complex.

Evidence that we are dealing with a single FMO complex stems from the following observations: First, for 320 pM concentration and a diffraction-limited excitation volume for the laser spot this yields a probability of 0.15 for finding a single FMO within this spot. Given that the FMO complexes have no tendency to aggregate, the probability to find two FMOs within the same volume is already less than 0.02. Based on these numbers one expects an average distance between two FMO complexes of about 15 μm , which is consistent with the widefield images of the samples. Second, for a single complex the emission rate n corresponding to the number of emitted photons per second is given by $n = \frac{\sigma}{A} \cdot \frac{P}{h\nu} \cdot \Phi_F$. Here σ denotes the absorption cross section, A the area of the focal spot, P the incident laser power, $h\nu$ the photon energy and Φ_F the fluorescence quantum yield. Inserting, $\sigma = 5.88 \times 10^{-20} \text{ m}^2$ (calculated from the extinction coefficient = $154 \text{ mM}^{-1} \text{ cm}^{-1}$ ⁴⁰), $A = 6.36 \times 10^{-13} \text{ m}^2$, $P = 80 \text{ nW}$, $h\nu = 2.48 \times 10^{-19} \text{ J}$, and $\Phi_F = 0.95$, we expect a fluorescence emission rate of about 29,000 photons/s. From previous work we know that the collection efficiency of our low-temperature microscope is about 2%, which yields 600 counts/s for the detected rate from a single FMO, in agreement with our observations.

Figure 3 compares the low-temperature static optical spectra of FMO from *Cba. tepidum* for a bulk sample and three different single complexes. In accordance with absorption spectra in the literature, the fluorescence-excitation spectrum from an ensemble (Fig. 3a, red line) features four bands, here labelled 1–4 in the order of increasing wavelength. Those can be characterised by four Gaussians (for details see experimental section) that decrease in widths going from band 1 to band 4, see Table 1. The corresponding spectrum from the single FMO complexes (Fig. 3a, black lines) in shape resemble closely the ensemble spectrum, with the exception that band 4 shows a lower relative intensity with respect to the ensemble spectrum and is sometimes hard to detect. In order to verify whether the single complex spectrum is affected by the excitation intensity we repeated the experiment for one single complex lowering and raising the excitation intensity by a factor of 5 and 4 respectively. The lower excitation intensity corresponds to the lowest intensity for which we can record an excitation spectrum. For these different excitation intensities we did not find a significant variation of the spectral profile.

Figure 3b displays the low-temperature emission spectra from three other individual FMO complexes together with the emission spectrum from a bulk sample. The samples have been excited at 805 nm and the excitation intensity was 250 W/cm^2 . Again we find a close resemblance between the ensemble and the single complex

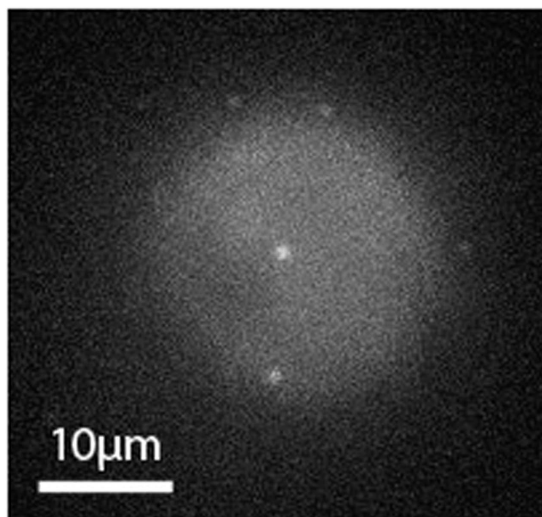


Figure 2. Widefield image of a sample with an FMO concentration of 320 pM. For excitation the output of the laser was wobbled between 800 nm and 809 nm at a rate of 0.3 s^{-1} . The excitation intensity was 175 W/cm^2 , and the acquisition time was 15 s.

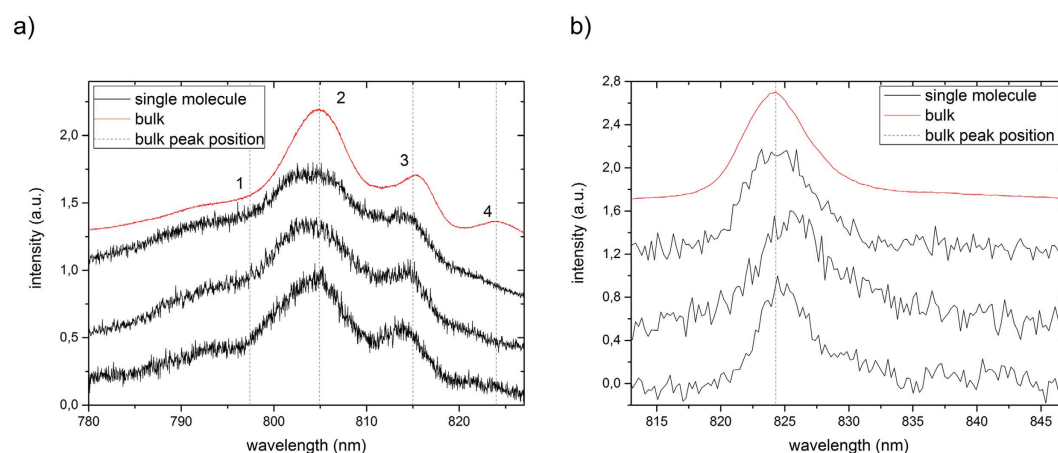


Figure 3. Low-temperature spectra (1.2 K) of the FMO complex from *Chlorobaculum tepidum*.

(a) Fluorescence-excitation spectra of FMO from an ensemble (red) and three different single complexes (black). All spectra were recorded with an excitation intensity of 12.5 W/cm^2 . The dashed vertical lines indicate the peak positions of the four Gaussians that have been used for fitting the bulk spectrum. For better comparison the spectra have been normalised, and are offset with respect to each other. The typical signal strength at the maximum amounts to 400 to 600 counts per second. Each spectrum corresponds to an average of 200 individual spectra recorded with an acquisition time of 19 s each. (b) Fluorescence emission spectra from an ensemble (red) and three different individual FMO complexes (black). The excitation wavelength was 805 nm, and the excitation intensity was 250 W/cm^2 . For better comparison the spectra have been normalised, and are offset with respect to each other. Each spectrum corresponds to an average of 400 individual spectra recorded with an acquisition time of 5 s each.

spectra. The emission spectrum features a single band that can be characterised by a Gaussian shape with a width of about 80 cm^{-1} (FWHM). For comparison, with one FMO complex we also recorded an emission spectrum with a reduced excitation intensity of 50 W/cm^2 . This resulted in an extremely noisy emission spectrum without significant change of its width. Details concerning the emission spectra are summarised in Table 2.

Discussion

The general finding is that there is no observable significant difference in the peak positions found in the linear optical spectra recorded for an ensemble or individual FMO complexes. We do, however, find a slight variation in the widths of these bands between individual complexes. Heating effects by the laser as observed for example in ref. 41 in the past as an origin of the line broadening can be excluded. In their work the samples were excited with pulses of 2 ps duration at 1 mJ/cm^2 corresponding to a peak excitation intensity of 500 MW/cm^2 .

		Bulk	complex 1	complex 2	complex 3
Peak 1	position (nm)	797.4 ± 0.5	795 ± 2	794 ± 2	796 ± 2
	FWHM (cm ⁻¹)	355 ± 8	490 ± 15	260 ± 15	316 ± 15
Peak 2	position (nm)	804.9 ± 0.5	804.8 ± 0.5	804.6 ± 0.5	804.5 ± 0.5
	FWHM (cm ⁻¹)	99 ± 2	160 ± 5	152 ± 5	133 ± 5
Peak 3	position (nm)	814.9 ± 0.5	814.5 ± 0.5	814.2 ± 0.5	814.0 ± 0.5
	FWHM (cm ⁻¹)	71 ± 2	88 ± 5	100 ± 5	75 ± 5
Peak 4	position (nm)	824.0 ± 0.5			
	FWHM (cm ⁻¹)	52 ± 2			

Table 1. Spectral positions and bandwidths observed in the fluorescence-excitation spectra displayed in Fig. 3a. The spectral peaks have been numbered 1–4 in the order of increasing wavelength. For the single-complex spectra a meaningful analysis of the positions and widths of peak 4 is prevented by the low signal-to-noise ratio.

	Bulk	complex 4	complex 5	complex 6
position (nm)	824.4 ± 0.6	824.6 ± 0.8	826.0 ± 0.8	825.1 ± 0.8
FWHM (cm ⁻¹)	97.0 ± 1.5	99 ± 7	114 ± 7	72 ± 7

Table 2. Spectral positions and bandwidths observed in the emission spectra displayed in Fig. 3b.

The concomitant rise in temperature was then estimated to be about 10 K if the energy is distributed over a group of seven BChl *a* molecules and 500 K if the energy is concentrated, for example by annihilation processes, on a single BChl *a* molecule. Here the highest excitation intensity used is 250 W/cm², i.e. more than 6 orders of magnitude lower than in the previous study. The widths of these bands, as predicted by theoretical work⁴², are determined mainly by the relaxation of the exciton states and/or fast fluctuations of the site energies of the individual pigments. Superimposed on this major contribution to the line-widths is a small additional broadening due to the heterogeneities (static and/or dynamic) between the complexes. Relaxation dynamics as the dominant origin of the line-widths is in line with the observation that the widths of the absorption bands decrease upon going from high- to low energy in the spectral positions of these bands. This is because the higher energy exciton states have more relaxation pathways available, which therefore accelerates their decay with respect to the low energy exciton states. The observed bandwidths of 70–400 cm⁻¹ provide a lower limit for the relaxation times of 10 and 70 fs, which is in agreement with previous work^{43,44}. Similarly, previous theoretical work predicts fluctuations of the site energies of the individual pigments on the same time scale^{42,45,46}. Obtaining a spectrum from a single FMO complex requires that the signal is collected for a finite dwell time, here 10 ms, in order to achieve a sufficient signal-to-noise ratio. Hence, a spectrum recorded from an individual FMO complex corresponds to a time-averaged spectrum, where the period over which the signal has been accumulated is very long with respect to the (predicted) time scale of the energetic fluctuations. In other words: A single FMO complex behaves ergodically on the experimentally accessible time scale. However, this single molecule study does reveal an important basic property of the internal energy transfer relaxation processes that take place within FMO, and it clarifies what kind of experiments with single FMO complexes can be done at all and what the relevant timescales are.

Methods

Sample preparation. FMO was prepared from cells of the thermophilic green sulphur bacterium *Cba. tepidum*. A strain containing a His-tagged version of the green sulphur bacterium *Cba. tepidum* reaction centre (RC) was a gift from the group of Prof. Oh-Oka⁴⁷. *Cba. tepidum* was grown anaerobically in the light in modified Pfennig's medium⁴⁸, known as liquid CL media^{47,49}. Glycerol stocks of the *Cba. tepidum* cells were inoculated in CL media and grown in 1.2 L air tight bottles that were allowed to go anaerobic by leaving them in the dark for 12 hours overnight. Cells were then grown at 43 °C under illumination for 2–3 days, at a light intensity of approximately 30 μmol photons m⁻² s⁻¹, before harvesting by centrifugation (12,000 × g). The RC was purified as described in ref. 50 with the exception of the cell breakage, which was carried out by using a cell disrupter (25 psi). At the stage of nickel affinity chromatography a high salt wash with >500 mM NaCl elutes FMO. This dilute solution of FMO was then purified by ion exchange chromatography on Whatman DE52 cellulose. FMO eluted between 90–200 mM NaCl. It was then concentrated by ultrafiltration with a 50 kDa MWCO membrane (Minicon concentrator, Millipore) and further purified, in 20 mM Tris HCl, pH8, by size exclusion chromatography on a Sepharose S-200 (GE Healthcare) column⁵¹.

The FMO complexes were stored in a 20 mM Tris A buffer (pH 8.0 at room temperature) at –80 °C until they were used. For single-molecule measurements the samples were diluted within three steps in a 20 mM Tris A buffer to a concentration of 320 pM. In the last dilution step the pre-diluted sample was mixed with the same amount of pure glycerol to avoid crystallisation at low temperatures. A small drop (0.5 μl) of this solution was seeped onto a cleaned SiO₂ plate before the drop was covered with a microscope coverslip to get a flat sample surface. The prepared sample was put immediately into the cryostat, pre-filled with liquid nitrogen to flash freeze the sample before it was cooled down to 1.2 K. For bulk measurements the sample preparation was identical, except for the FMO concentration which was 800 μM.

Fluorescence-Excitation Spectroscopy. For the fluorescence-excitation experiments the samples are excited with the output from a titan-sapphire-laser (3900 S, Spectra Physics) that was pumped by a Nd:YVO₄ laser (Millenia Vs, Spectra Physics). The wavelength is varied between 770 nm and 827 nm by turning the birefringent filter with a stepper motor (Actuator 850 F, Motion Controller MM4005, Newport). The accuracy and the reproducibility of the wavelength variation is 1 cm⁻¹, as verified with a wavemeter (WaveMaster, Coherent). The excitation light passes a home-build fluorescence microscope and is focussed by an objective (Mikrothek, NA = 0.85) that is mounted inside the cryostat to a diffraction-limited spot of 0.9 μm in diameter. For the selection of an individual complex the microscope is operated in wide-field mode and the laser wavelength is wobbled between 800 nm and 809 nm. The emission from the sample is collected by the same objective, transmitted through two band-pass filters (LP830, AHF Analysetechnik) for suppressing residual laser light, and detected with a CCD-camera (iKon, Andor). From the widefield image (Fig. 2) a spatially well separated FMO complex is selected and the microscope is switched to confocal mode. A similar setup is described in great detail in refs 52,53. The signal passes the bandpass filters mentioned above and is registered with a single-photon counting avalanche photodiode (APD) (SPCM-AQR-16, Perkin Elmer). During acquisition of the excitation spectra, the intensity of the laser is recorded as a function of the wavelength using a powermeter (LaserMate-Q, Coherent), and all spectra are corrected for variations of the laser intensity.

Since in the literature the FMO spectra are mostly given on a wavelength scale we have chosen to display the spectra in Fig. 3 on a wavelength scale. However, for fitting the data we used an energy equivalent axis (wavenumbers) on the abscissa and took the required intensity correction into account for the conversion of the spectral peak positions from wavenumbers to the corresponding wavelengths.

Emission Spectroscopy. For emission spectroscopy the sample is excited at 805 nm, and the light from the sample is guided through a spectrometer (Acton 250, Princeton Instruments) that accommodates a mirror and a grating (300 lines/mm, blaze 1000 nm). For selecting a single FMO complex the setup is operated in wide-field mode and the emission from the sample is directed onto a CCD-camera (iKon, Andor) via the mirror in the spectrometer. From the widefield image a single FMO complex is selected, the microscope is switched to confocal mode as before, and the mirror in the spectrometer is replaced by the grating. For registering the spectra with the CCD, 16 pixel were binned vertically to reduce the read-out noise. For a single complex 400 spectra were recorded successively with a dwell time of 5 s, which yields a total acquisition time of 30 minutes. Since the individual spectra did not feature any fluctuations on these time scales we display the sum spectrum in Fig. 3b.

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Author Contributions

A.L. conducted the experiments, analysed the data and helped to write the paper. K.A. prepared the samples, R.J.C. supervised sample preparations and designed research. J.K. supervised the experiments, designed research and wrote the paper. All authors reviewed the manuscript.

Additional Information

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